

Properties of High-Intensity EUV & Soft-X Radiation Plasma Sources

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fire
Fluid, Ions and Radiation Ensemble
in Integrated Plasma Modelling



Sources for EUV Lithography

Diffraction restricts
the resolution

$$r \geq k_1 \frac{\lambda}{NA}$$

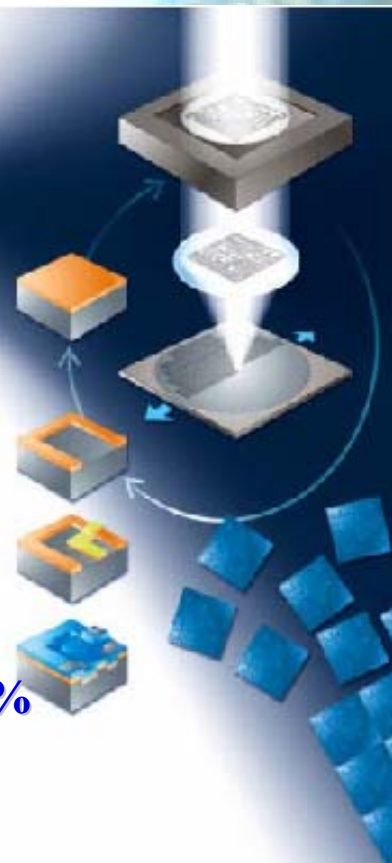
NOW
EUV for HVM
beyond 16 nm

$\lambda \Rightarrow 13.5\text{nm} \Rightarrow 6.X\text{nm}$
 $(h\nu=92\text{eV} \Rightarrow 185\text{eV})$

$\delta\lambda/\lambda \Rightarrow 2\%$

Nano-Age World

The optics is made of
multi-layer mirrors
with reflection efficiency ~70%



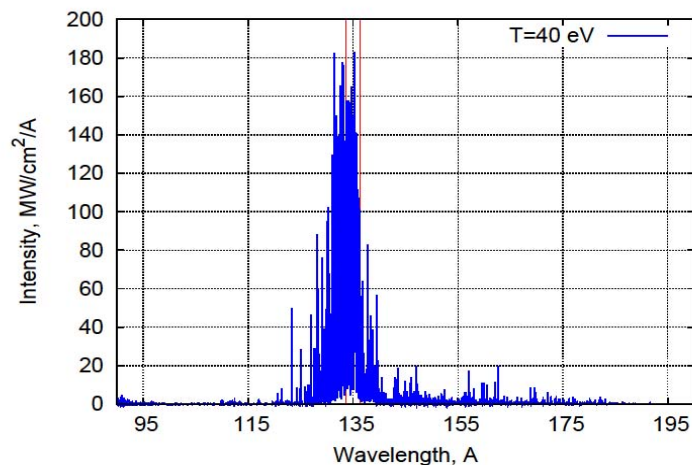
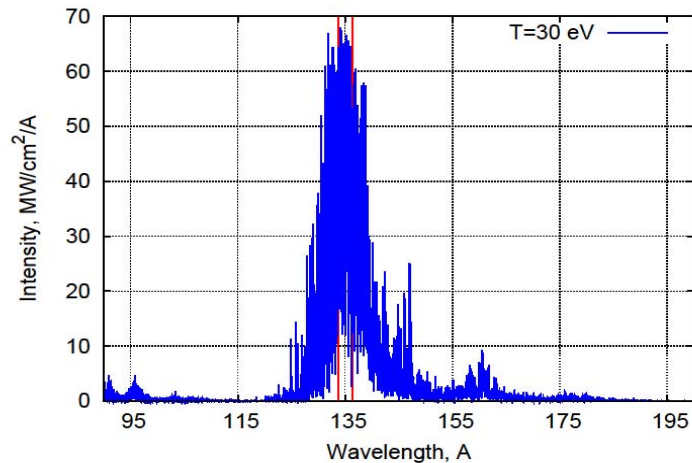
- For HVM: $\gg 200\text{ W}$ of in-band power at IF within $< 3\text{mm}^2\text{sr}$ etendue
- For mask inspections ABI→AIMS→APMI : $30 \rightarrow >100\text{ W/mm}^2\cdot\text{sr}$

Sn, Xe... High Energy Density plasma
($T_e=20\text{-}40\text{eV}$) radiates in EUV range
LPP & DPP



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EUV Brightness Limit for EUV Source



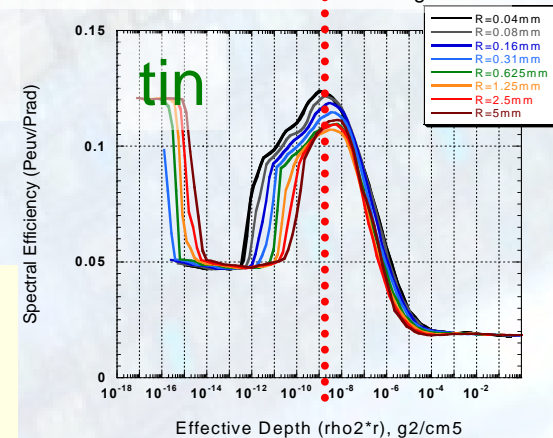
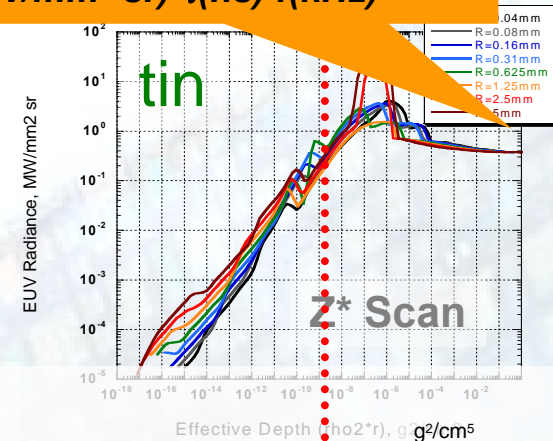
Detailed spectra from tin plasma with radius $R=100 \mu\text{m}$ and $n_e=10^{19} \text{ cm}^{-3}$

**Spherical
model of tin
plasma EUV
source**

RMHD scan for tin
plasma optimized
by radius, temperature
and density [AL10]

**The radiation self-
absorption limits the
in-band EUV radiance
from the plasma, and
the etendue
constraint limits the
usable power at IF of
a conventional single
unit EUV source**

$$L \approx 1.1 (W/mm^2 \text{ sr}) \cdot \tau(ns) \cdot f(kHz)$$



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EUV Brightness Limit at Higher Temperature

LTE

- The intensity upper Planckian limit of a single spherical optically thick plasma source in $\Delta\lambda/\lambda=2\%$ band around $\lambda=13.5\text{nm}$

$$I = \frac{2hc^2}{\lambda^4} \frac{\Delta\lambda/\lambda}{e^{\frac{hc}{\lambda T}} - 1} \approx \frac{72}{e^{\frac{92}{T(\text{eV})}} - 1} (\text{MW} / \text{mm}^2 \text{sr})$$

- Source with pulse duration τ and repetition rate f yields the time-average radiance $L = I \cdot (\tau f)$

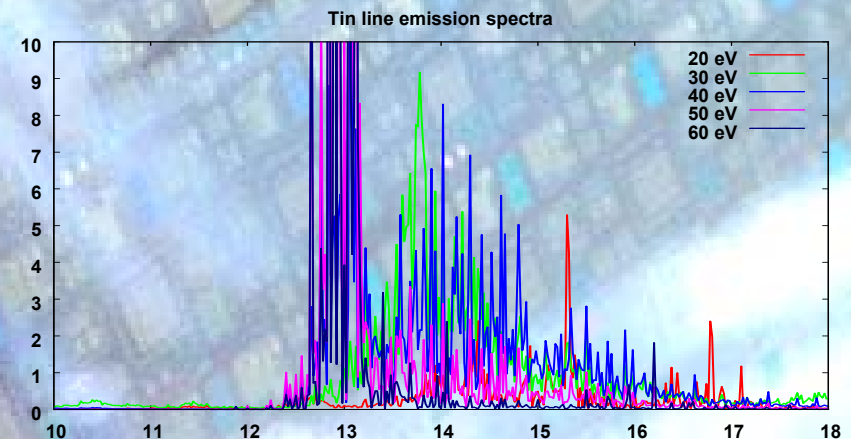
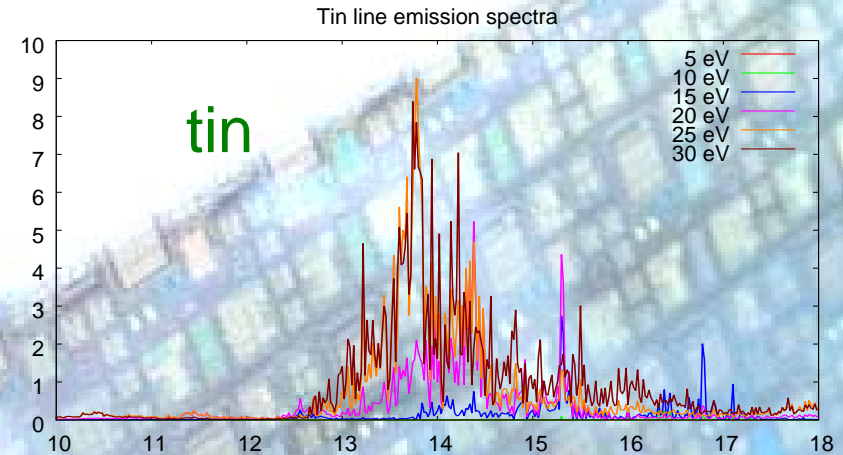
- The spectral efficiency has the maximum at $T \approx 22\text{eV}$

$$L \approx 1.1 (\text{W}/\text{mm}^2 \text{sr}) \cdot \tau(\text{ns}) \cdot f(\text{kHz})$$

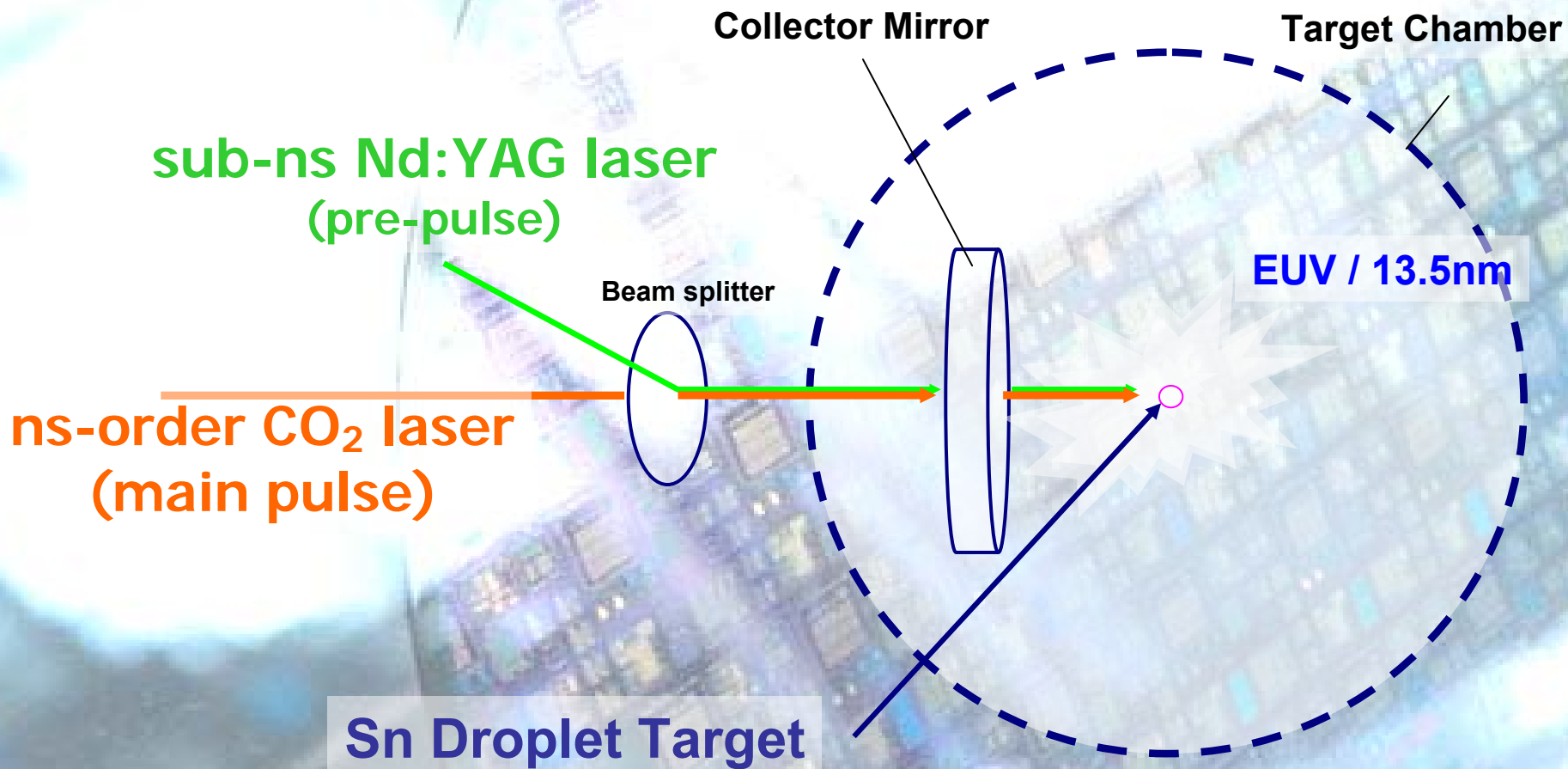
- For instance, at $\tau = 20\text{ns}$

$$L = 22 (\text{W}/\text{mm}^2 \text{sr})/\text{kHz}.$$

Non-LTE

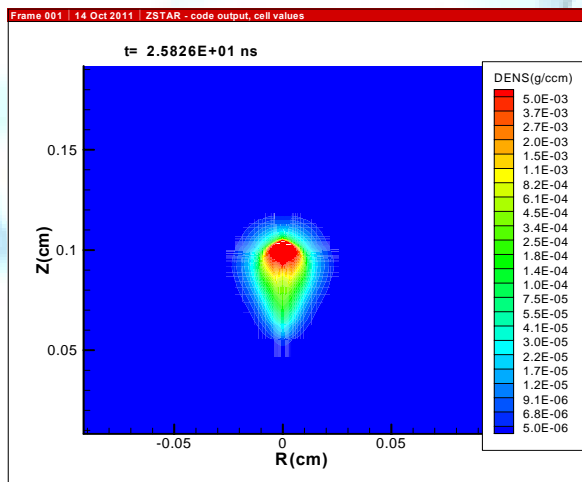


Combined Nd:YAG - CO₂ Laser System

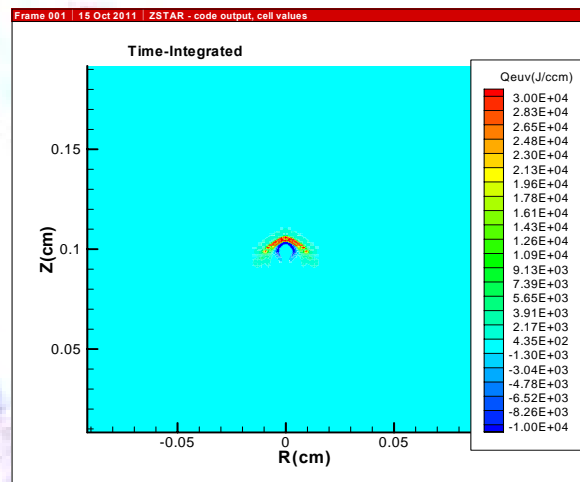


LPP EUV Source

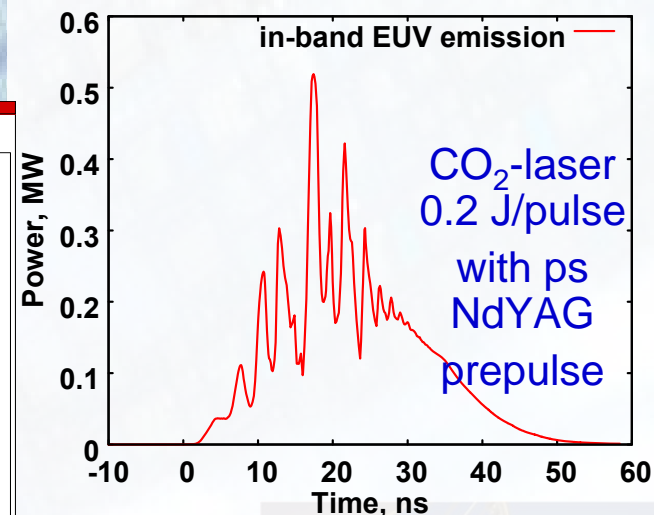
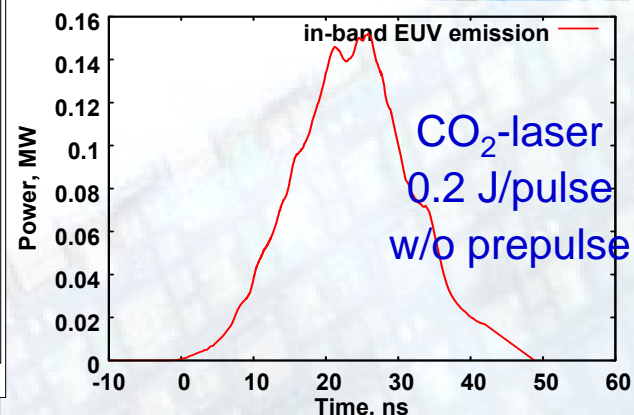
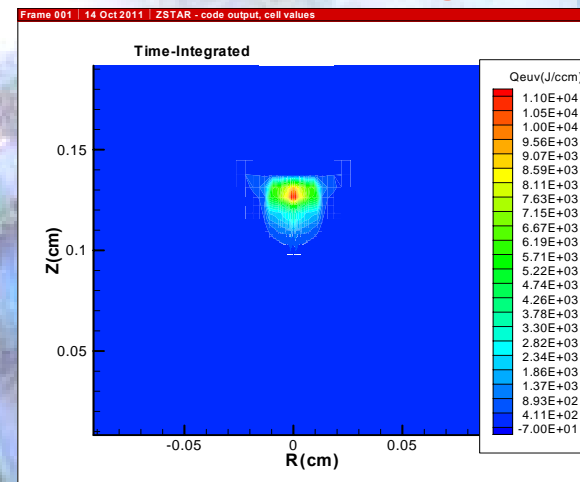
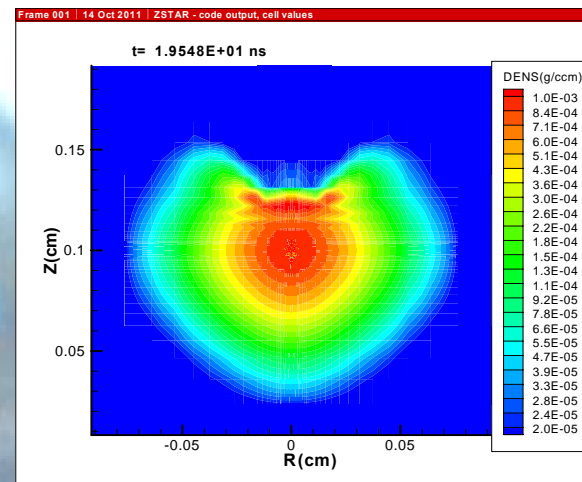
under CO₂- laser or combined pulse



Tin plasma density at
EUV maximum



Time-integrated
EUV source image



The maximum EUV brightness is up to 15 W/mm² sr kHz

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Conversion Efficiency of CO₂-laser

on pulse duration, with & w/out pre-pulses

Main pulse: CO₂-laser 0.1-0.8 J/pulse, 10,15,30,50ns fwhm, 200 μ m focal spot

Pre-pulse laser (if applied): Nd:YAG 5 mJ 1-10ns pulse, 40 μ m spot size

or Nd:YAG 6 mJ 10-100 ps pulse, 40 μ m spot size

Target:

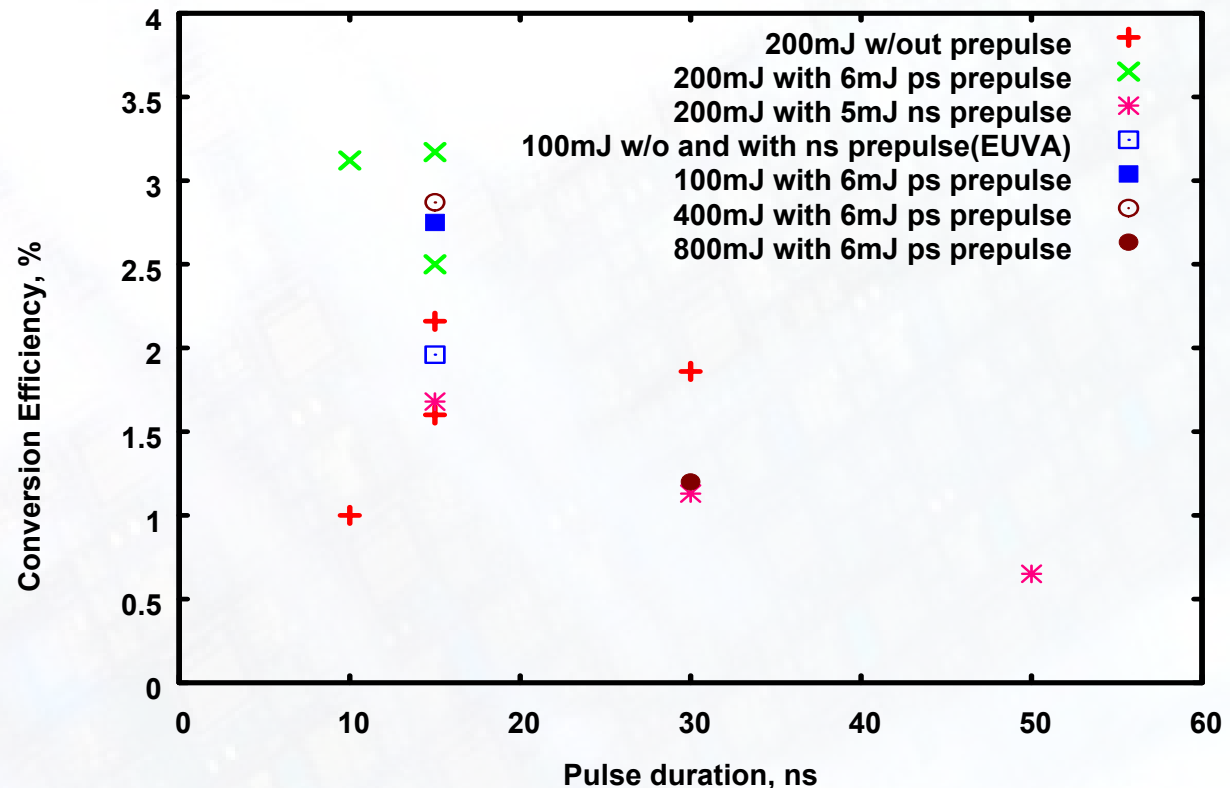
Liquid tin droplet of 40 μ m diameter

or
20 μ m for 100mJ (EUVA)

Conditions:

different focal positions;
different time delay between
pre- & main-pulse

**CE depends strongly on laser
intensity and target irradiation
conditions**

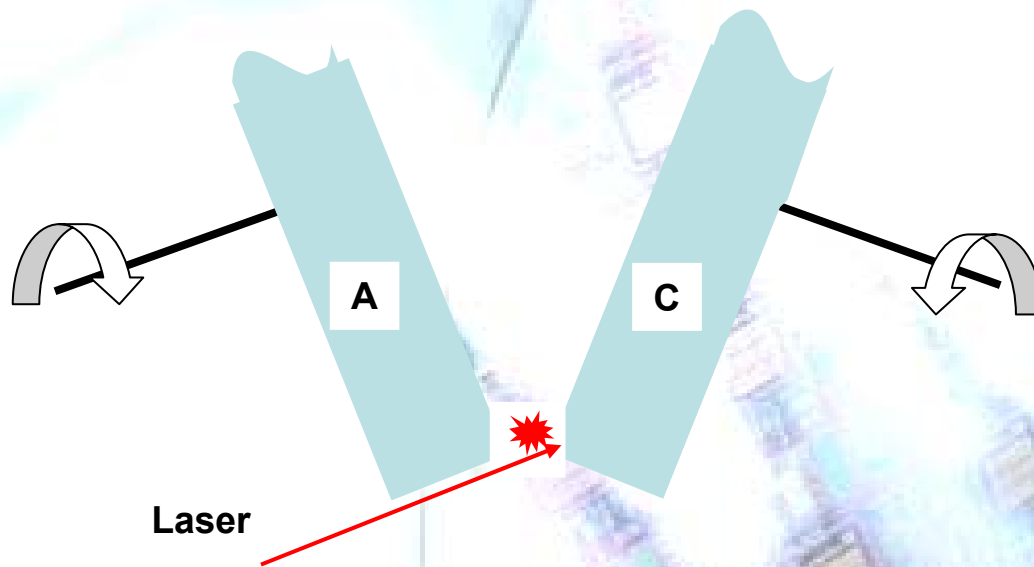


**CE maximum of 3% can be reached at laser energy (200mJ)
in a combined ps-ns pulse**

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Laser Assisted Vacuum Arc (LAVA-lamp)



High-current discharge between two rotating electrodes covered with a thin liquid Tin or Galinstan film is triggered by local laser ablation of electrode material.

Discharge

capacitance	0.4 μ F
inductance	19 nH
voltages	3 – 6 kV
energies	1.8 – 7.2 J
current	20 kA at 4.5 kV

Trigger laser:

wavelength	1064 nm
beam diameter	3 mm
focal lens	30 cm
energy	5 – 50 mJ
(varied by means of rotatable half-wave plate and polarizing beam splitter)	

Details are presented in the posters :

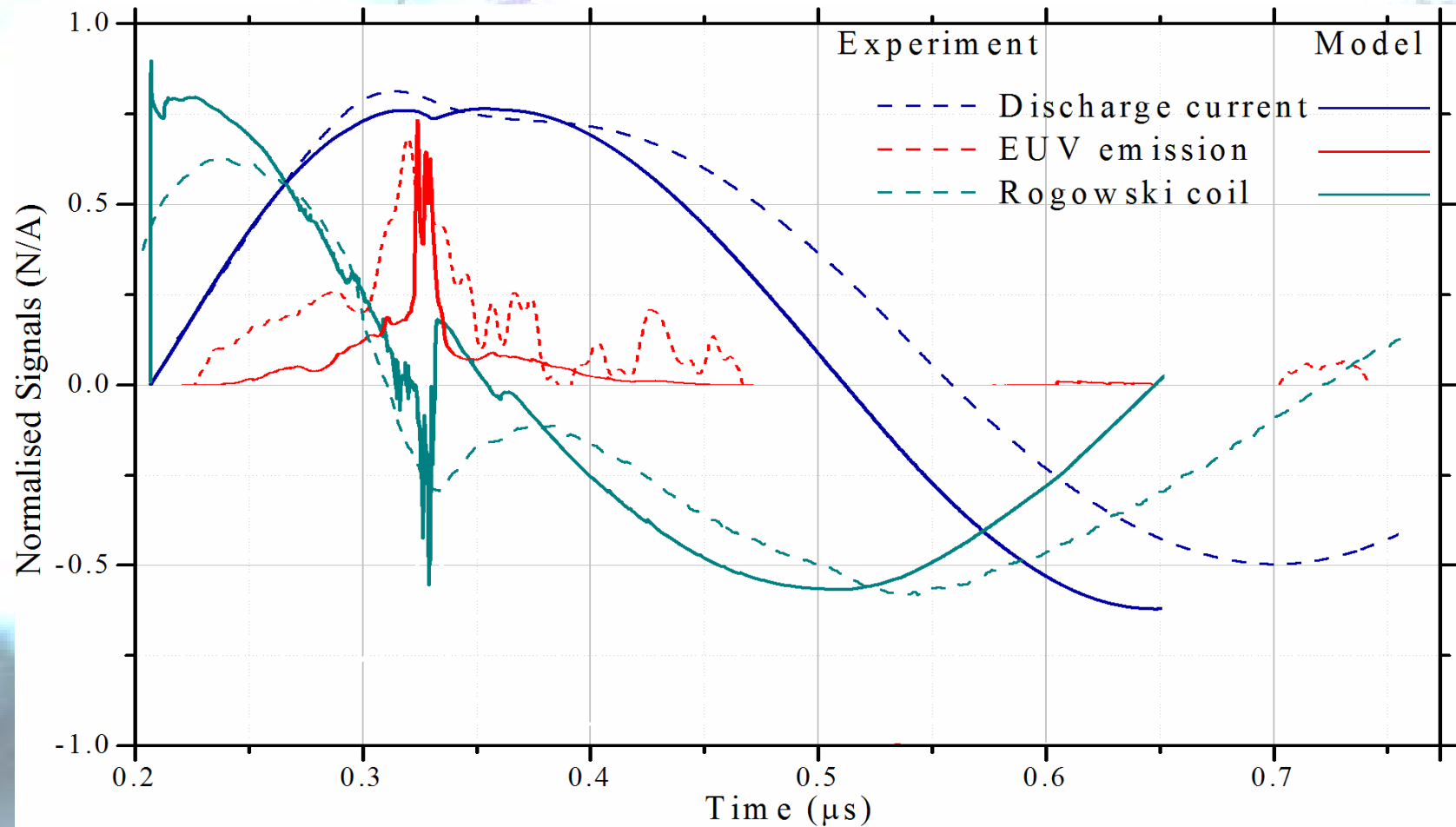
S26 V.S. Zakharov, Larissa Jushkin, S.V. Zakharov et al

S30 Isaac Tobin, Larissa Juschkin, Vasily S. Zakharov et al

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Comparison of measured and Z^* modelling discharge current and in-band EUV emission

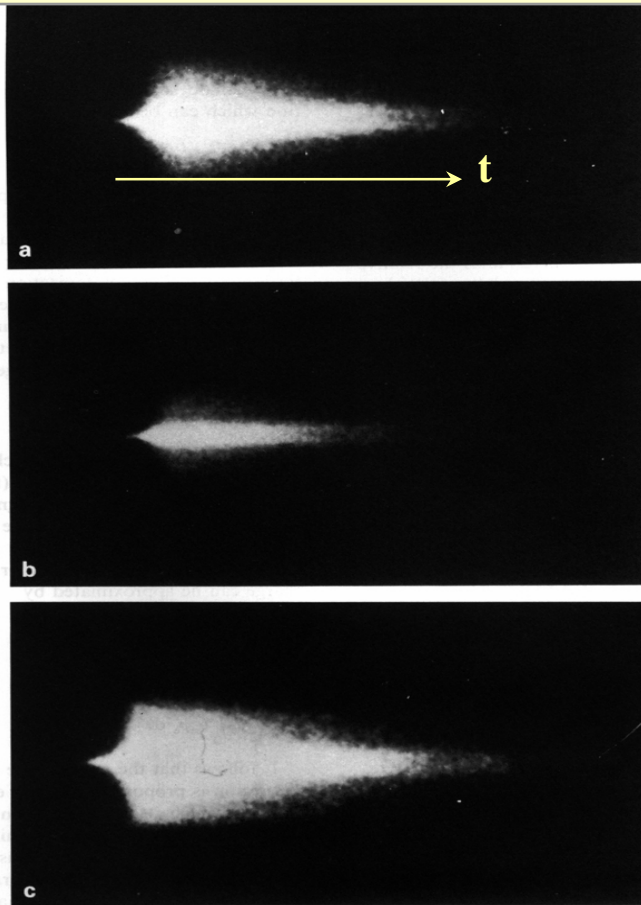


tin

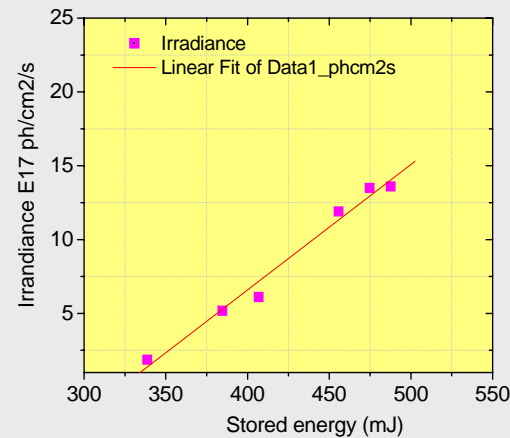
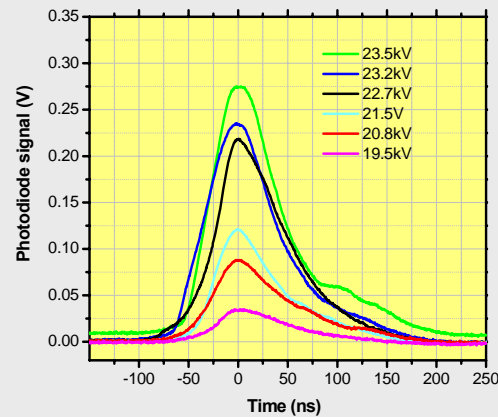
Capillary Discharge EUV Source

EXPERIMENT: discharge glow & EUV emission

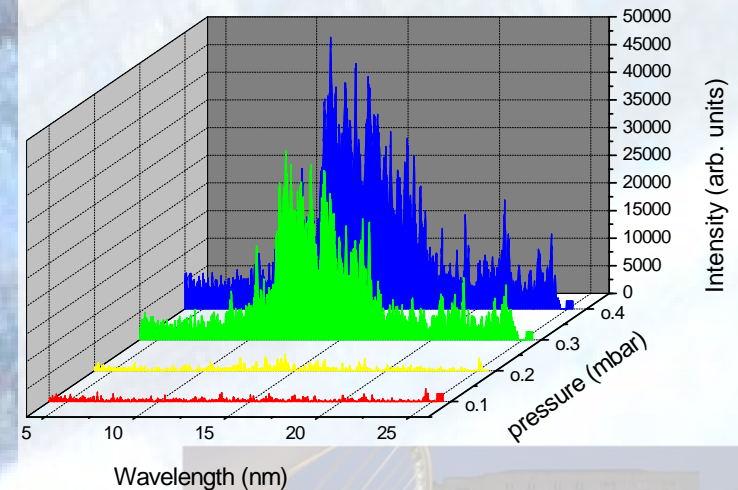
optical streak photograph



EUV emission



EUV spectrum



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High Brightness EUV Plasma Source

pulsed capillary discharge

Power source

Charge energy 0.2 – 0.5 J

Current 5 - 10 kA

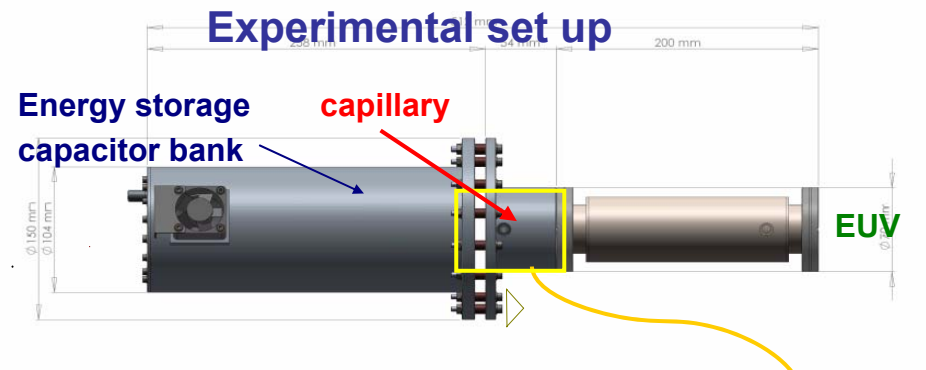
Pulse ~10-20 ns

Capillary dimension: \varnothing 1.6 mm
L = 12-18 mm

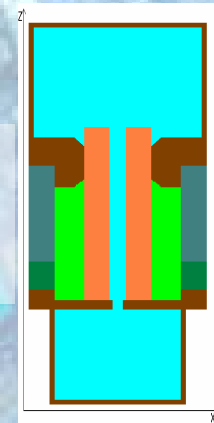
Various electrode geometries

Gas:

0.02-2 Torr gradient He;
Xe, N₂, Ar, Kr,, ... admixtures
(for narrow-band radiation source)



Example of
simulated
geometry



Capillary discharge dynamics & emission features:

E-beam, plasma channelling ($\epsilon \gg 1$)

Volumetric MHD compression (skin depth \gg plasma diameter)

Highly ionized ions (fast electrons)

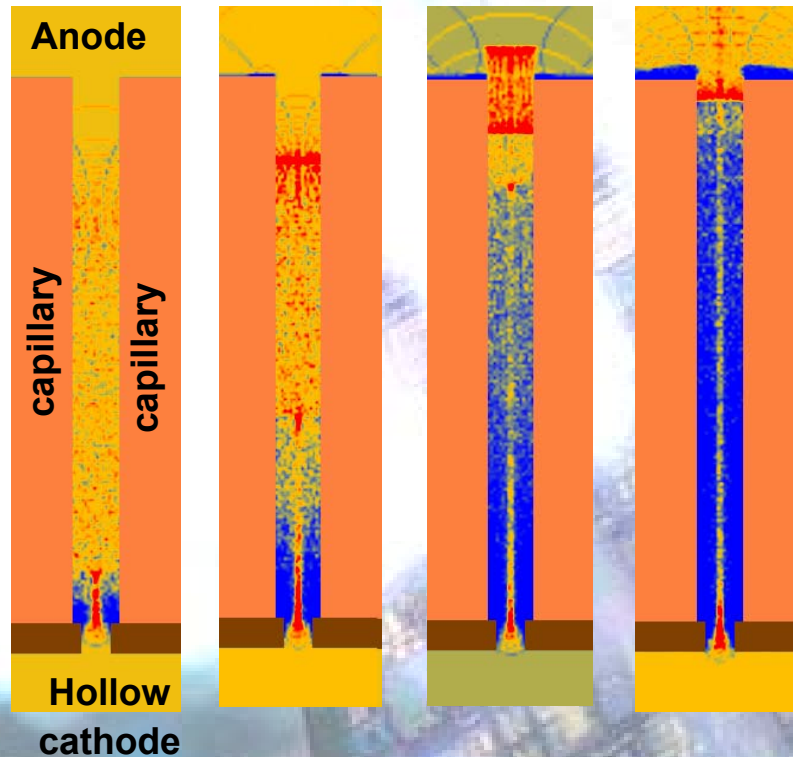
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Hollow-cathode Capillary Discharge

modelling: triggering by fast electrons

together with Markov M.B. et al, KIAM RAS



Electron beam in the HC capillary discharge

⇒ run-away electrons

⇒ electric field drops deeper into HC

⇒ e-beam concentration ($\varepsilon \gg 1$)

⇒ e-beam-gas ionization

⇒ ionization wave

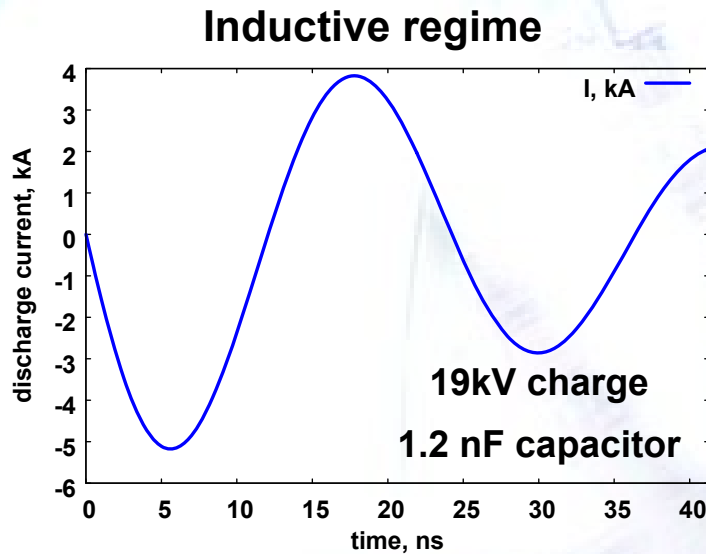
In the first few nanoseconds, run-away electrons from the hollow cathode generate a tight ionized channel ($< 200\mu\text{m}$ diameter) in the gas

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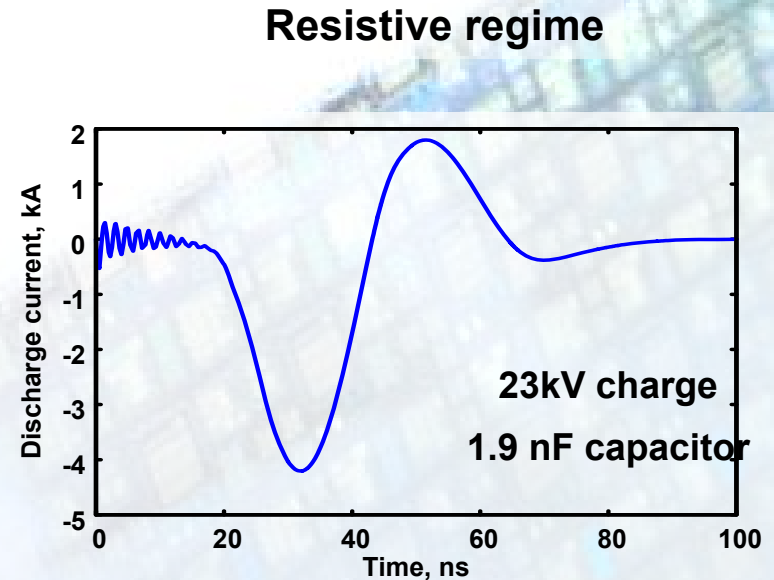
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Capillary Discharge EUV Source

Z*-code modelling: resistive regime



Nitrogen as
buffer gas



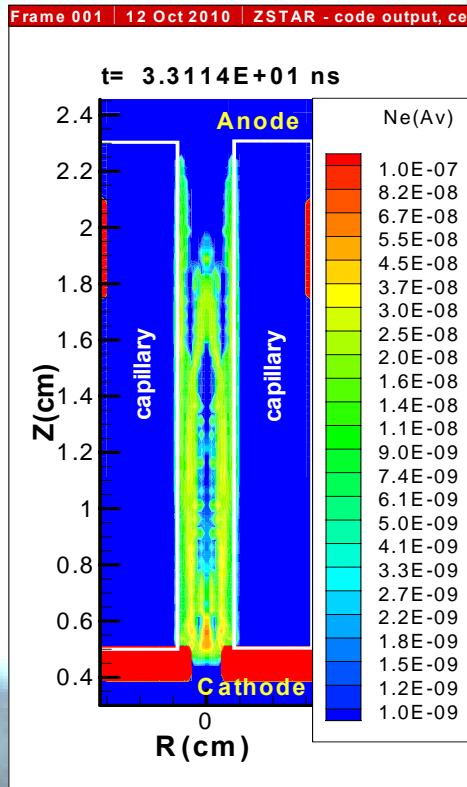
In the resistive regime of capillary discharge, the high joule dissipation in the tight conductive channel produced by hollow cathode electron beam creates an efficient mechanism of plasma heating and EUV or soft X-ray emission.

Also, fast electrons increase the ionization degree of heavy ion (Xe,...) plasma increasing eo ipso EUV yield.

Capillary Discharge EUV Source

dynamics & EUV emission

3D volumetric
compression

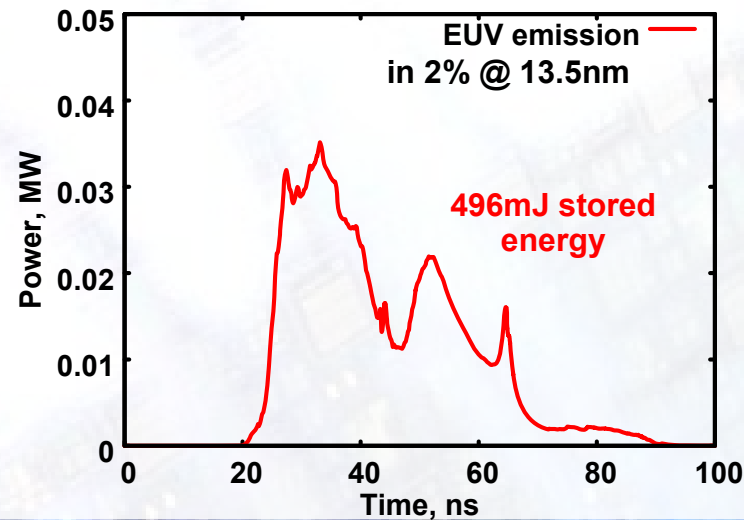


The traced along the axis, EUV
intensity at 13.5nm wavelength

15.3 W/eV mm² sr per kHz

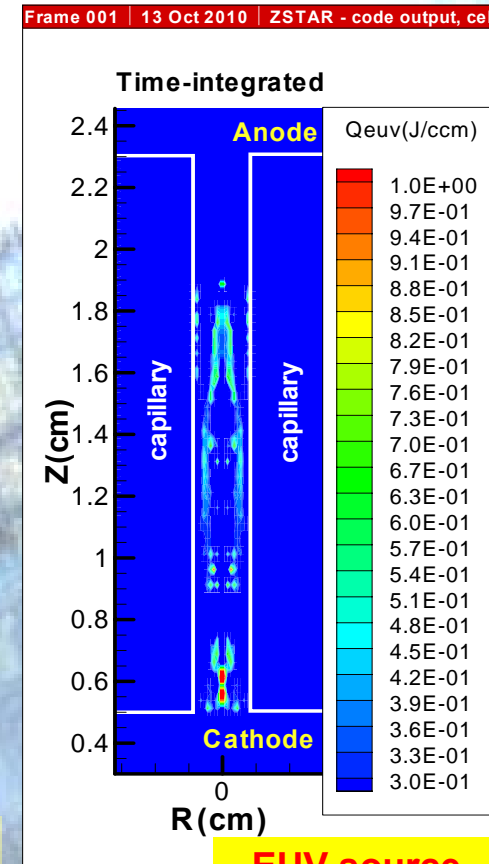
$$N_e = 2-3 \cdot 10^{17} \text{ cm}^{-3},$$

$$T_e = 25-40 \text{ eV}.$$



Calculated in-
band EUV
emission

0.885 W/kHz



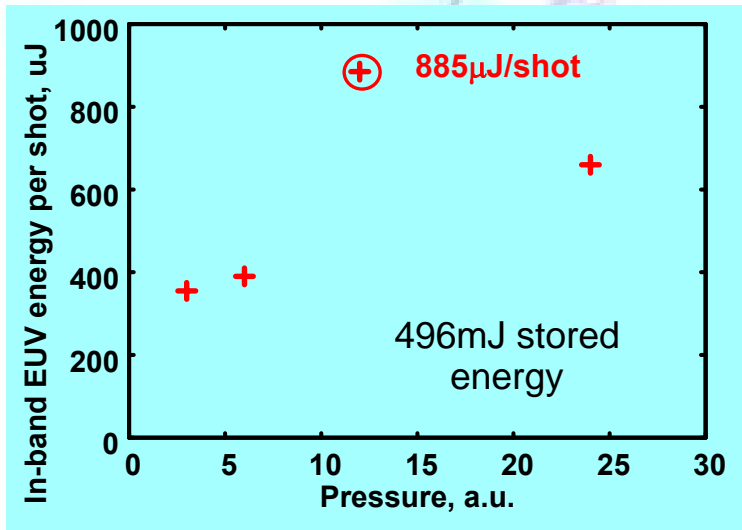
EUV source
cross-section

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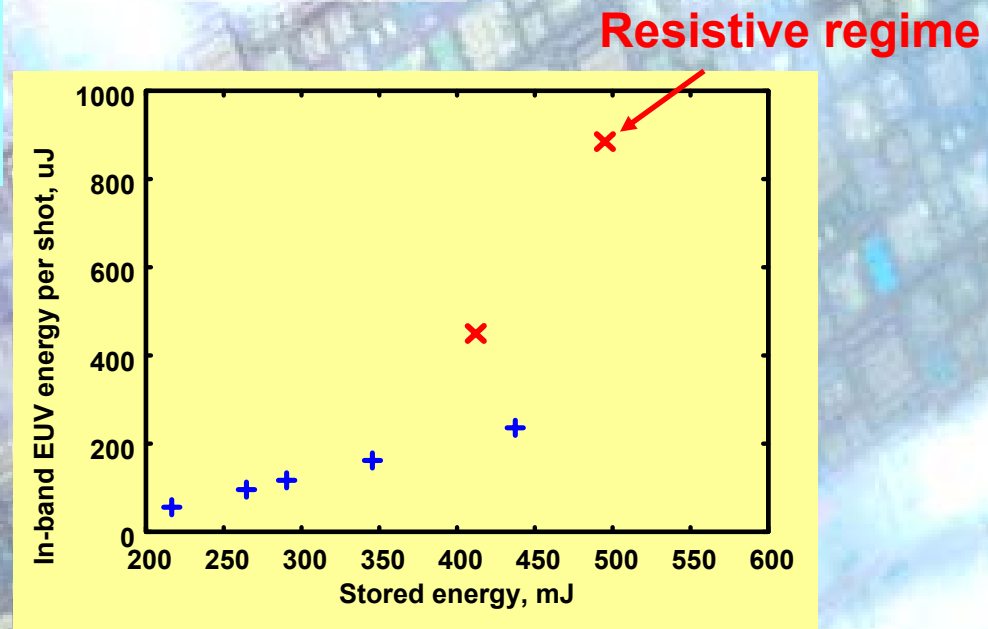
Capillary Discharge EUV Source

Z*-code modelling: source optimization



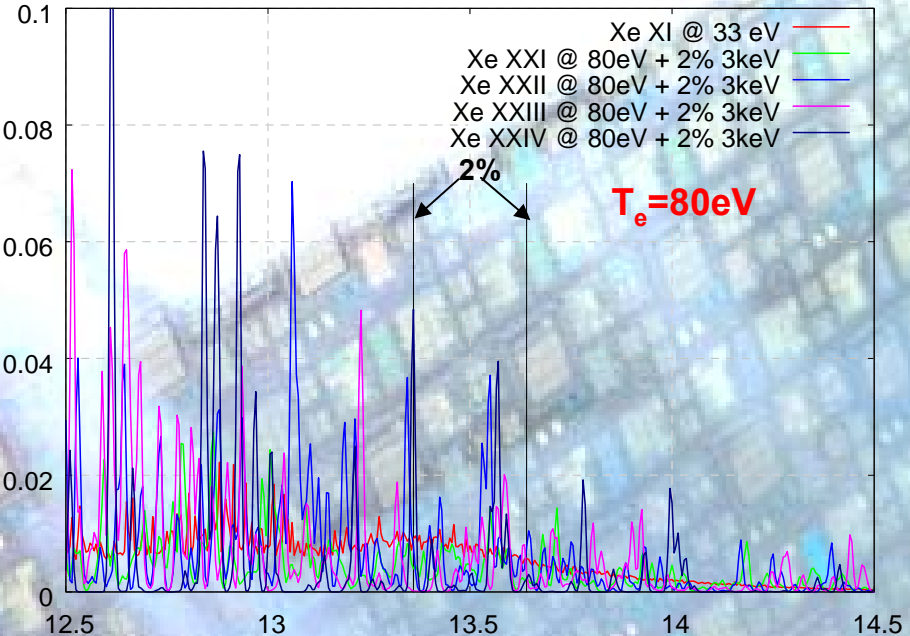
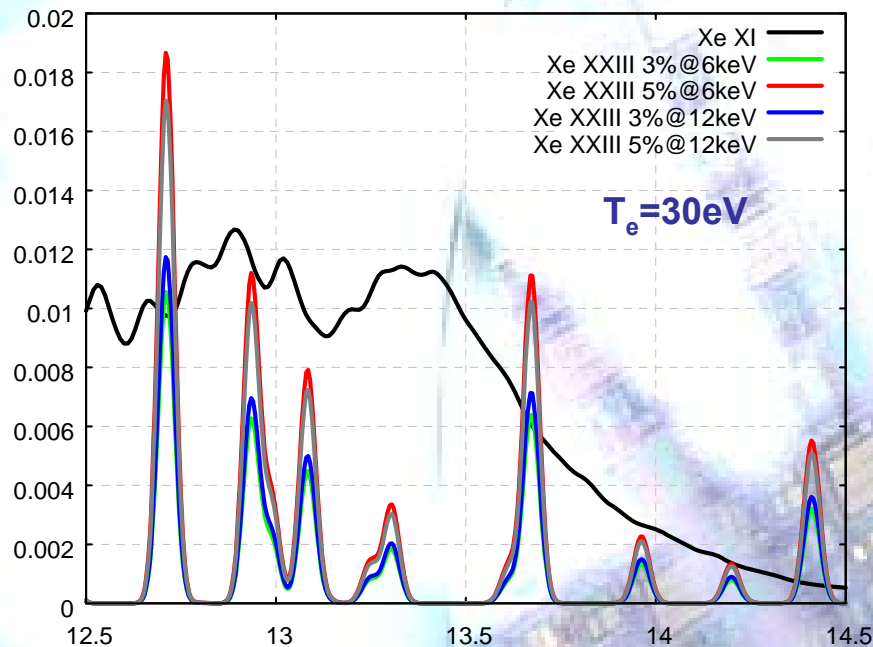
EUV source scan
by stored electrical
energy

Optimization
by gas mixture
pressure



EUV Emission of Highly Charged Xe Ions

- from plasma with fast electrons

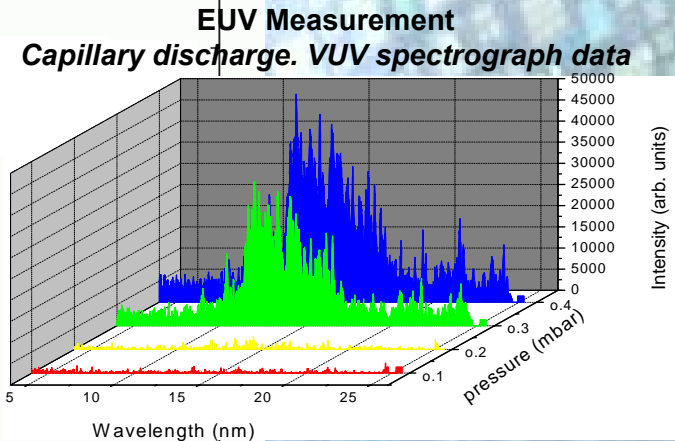
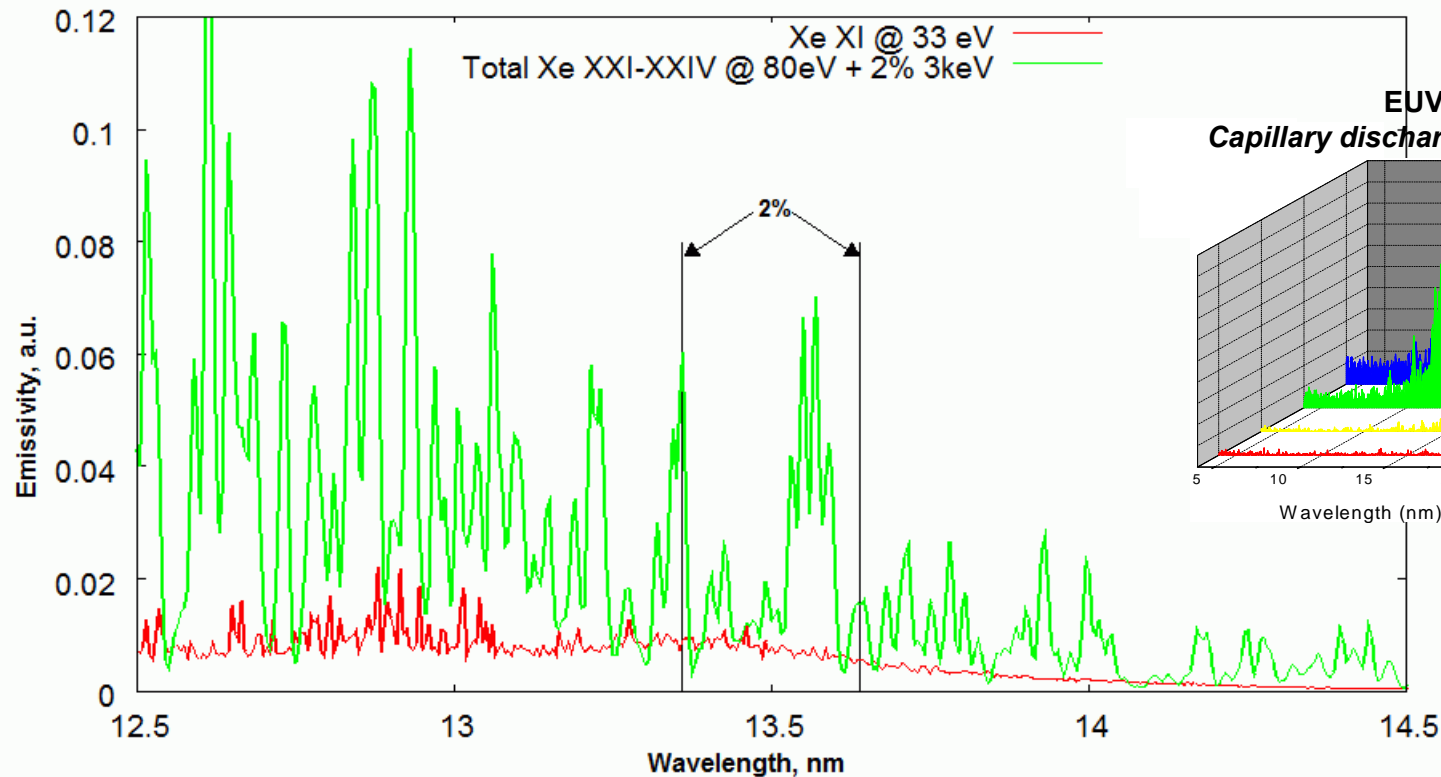


To produce the maximum EUV light power the double condition is required:

- + fast electrons have the energy of few keV to produce the highly charged ions
- + plasma has the temperature sufficient for the excitation of required transitions

EUV Emission of Highly Charged Xe Ions

- from e-beam triggered discharge plasma

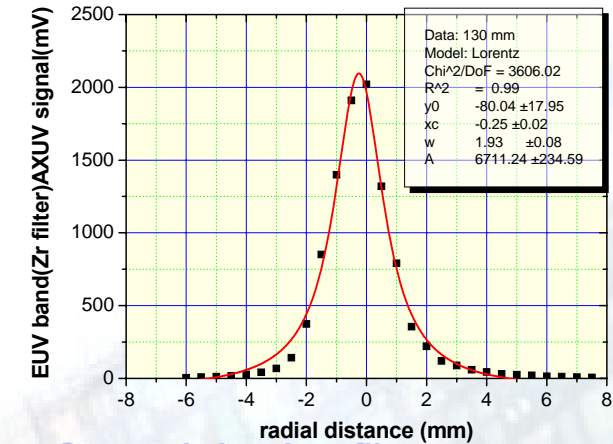
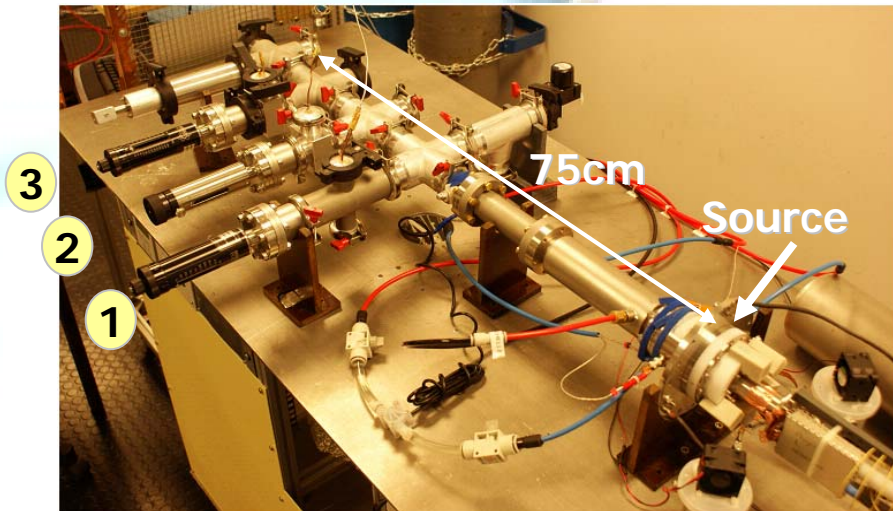


Bright EUV emission in 2% band at 13.5 nm can be achieved from highly charged xenon ions in plasma with small percentage of fast electrons

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Focusing Effect Observation



Scanned signal profile

EUV band (Zr filter) radiation beam profile at 130mm from collimator exit

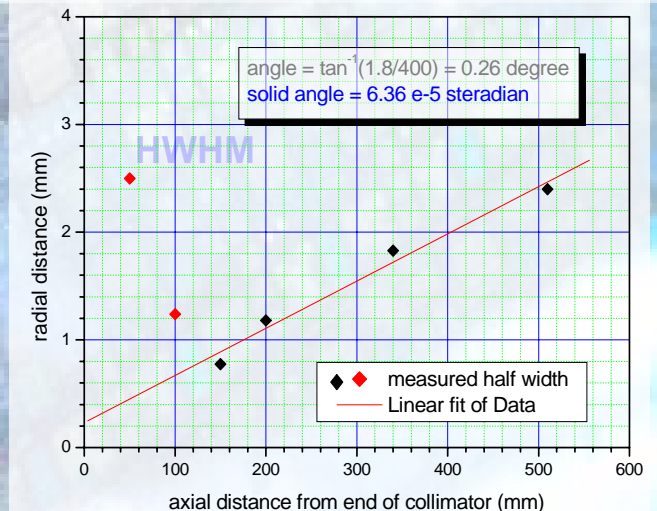
$$n^2 = 1 - \frac{\omega_e^2}{\omega^2} f_1(\omega)$$

$$\delta n = |1 - n| \ll 1;$$

$$\delta n \sim 0.01 \div 0.05 \text{ (in solid matter) and}$$

$$\delta n = 0.0000 \dots \text{ (in plasma) for EUV range}$$

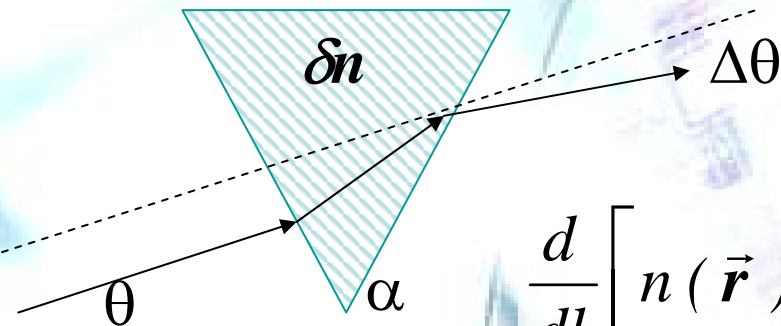
How it is possible in geometrical optics?
Know - How



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Wave-guiding Refractive Structure



$$N = \frac{\theta}{\Delta n \cdot \sin(\alpha)}$$

refractions are required

$$\frac{d}{dl} \left[n(\vec{r}) \frac{d\vec{r}}{dl} \right] = \vec{\nabla} n(\vec{r})$$

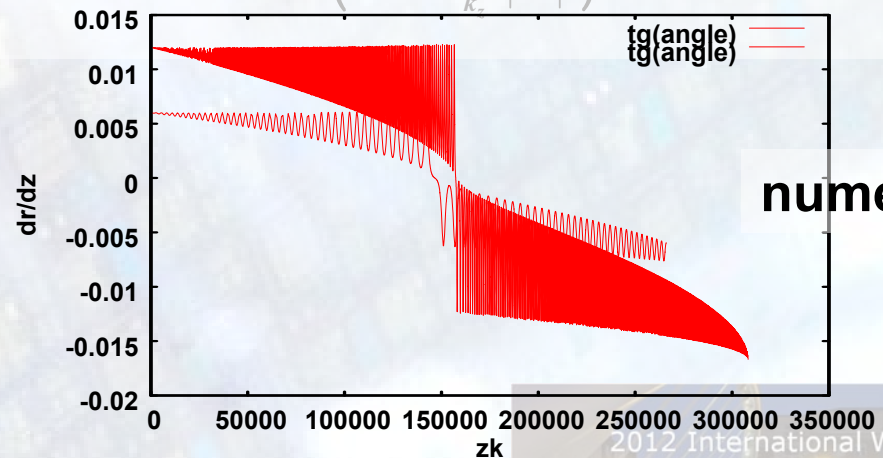
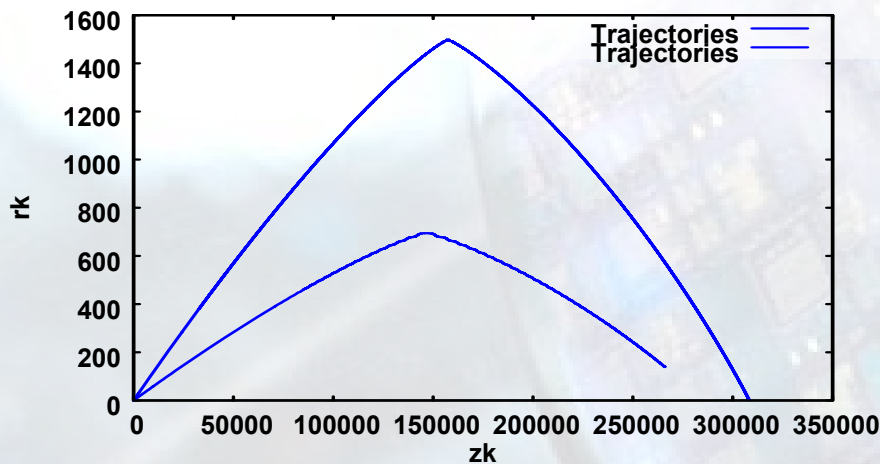
light trajectory equation

Refractive Structure: e-beam excited Kielwasser-waves, $k \leq r_D^{-1}$

Focussing :

$$\bar{\theta}(z) - \theta_0 \approx -0.25 \int \left(\frac{k_r |\Delta n|^2}{1 + 0.5 \frac{k_r^2}{k_z^2} |\Delta n|^2} \right) dz$$

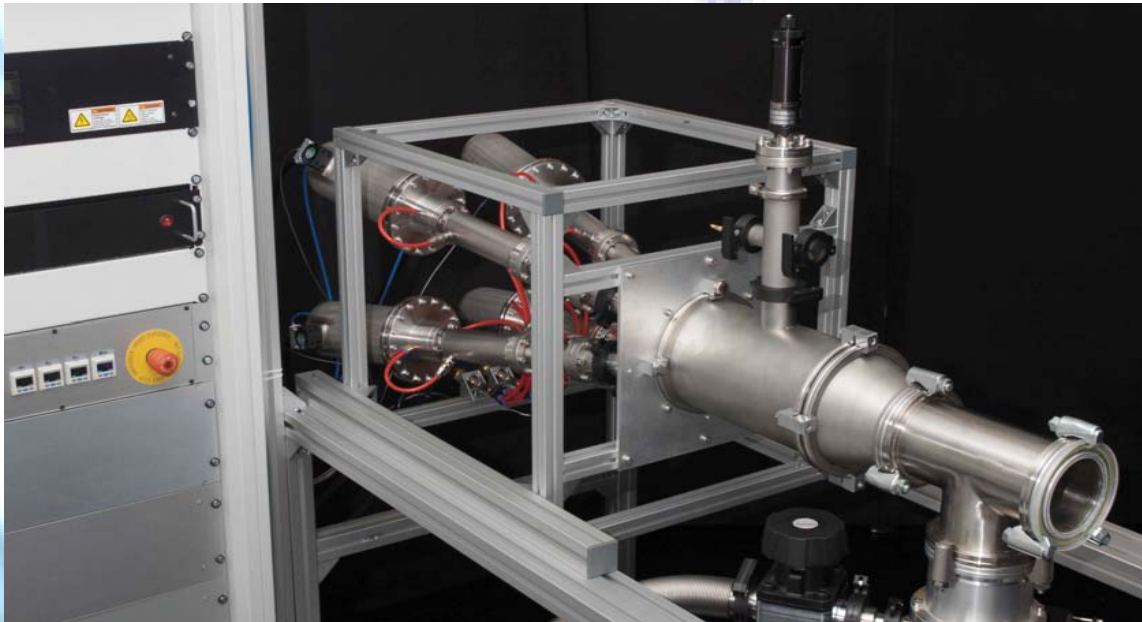
analytical



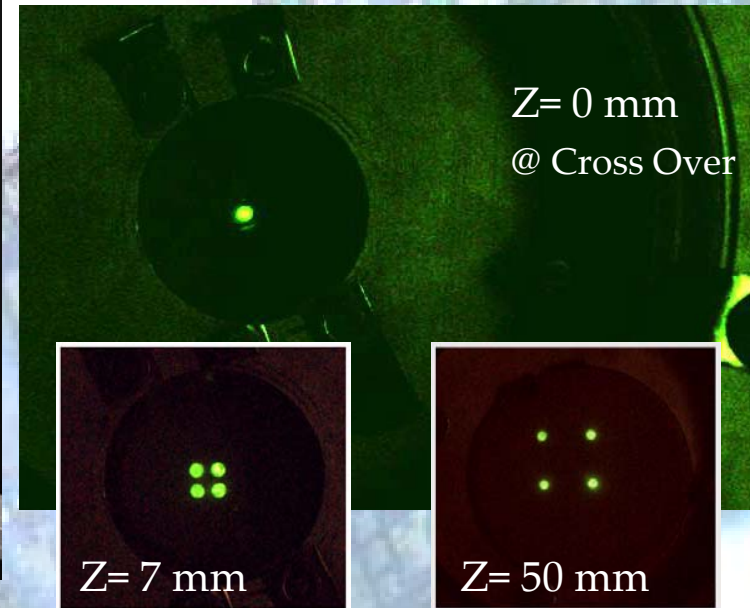
numerical

Multiplexer ⁴ :

- spatial multiplexing



4 sources operating individually
with common power control



All 4 sources aligned to a point
without use of any solid optical collector

Multiplexer 4: Optical Schematic

static combination of source beams to one

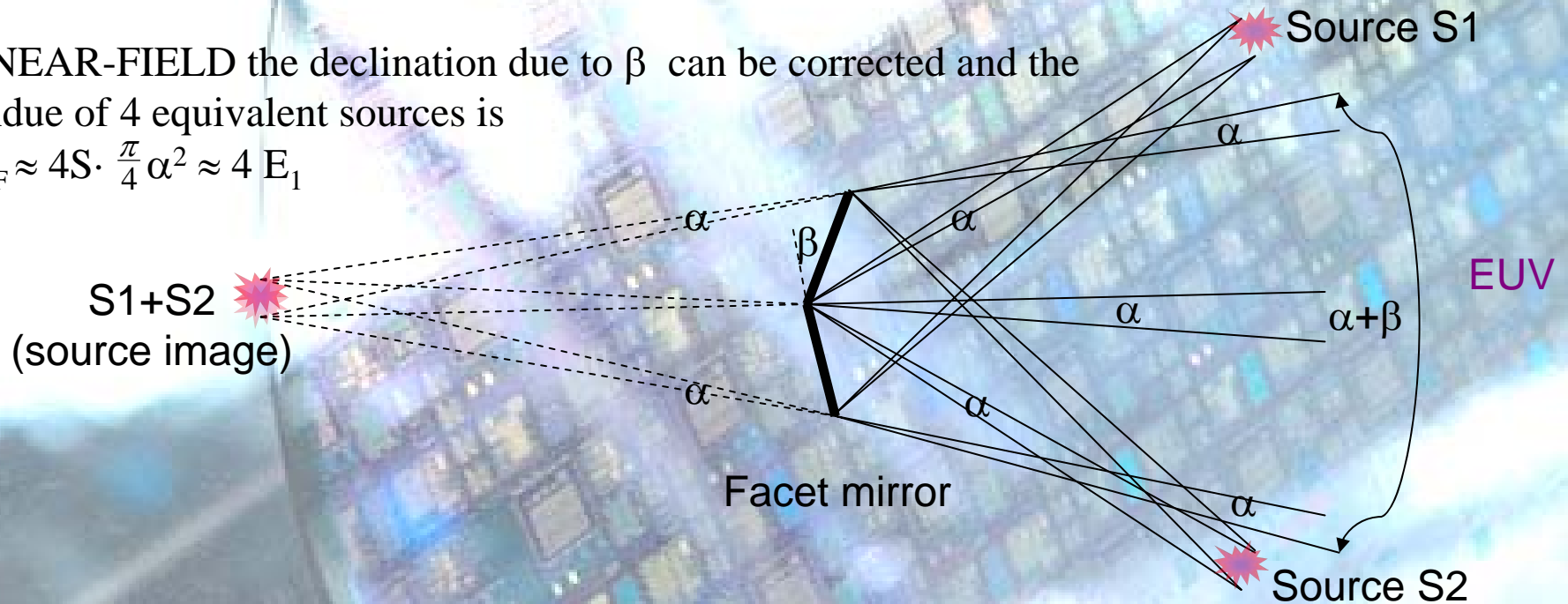
Etendue of a single source is $E_1 \approx S \cdot \alpha^2 \frac{\pi}{4}$

IN FAR-FIELD the etendue of 4 equivalent sources is

$$E_{4FF} \approx 4S \cdot \frac{\pi}{4} (\alpha + \beta)^2 \approx 16 E_1$$

IN NEAR-FIELD the declination due to β can be corrected and the etendue of 4 equivalent sources is

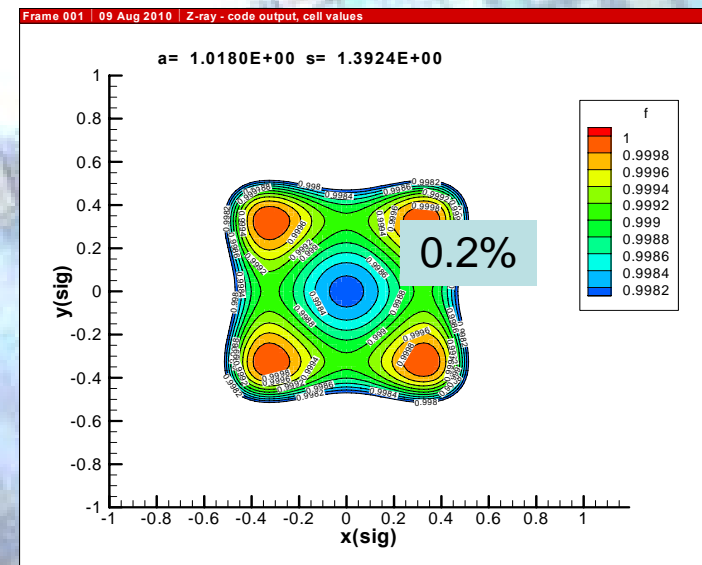
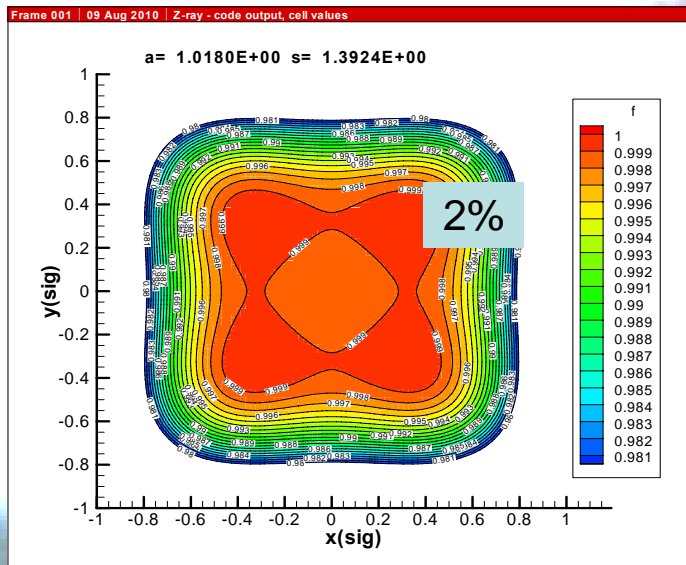
$$E_{4NF} \approx 4S \cdot \frac{\pi}{4} \alpha^2 \approx 4 E_1$$



Multiplexer⁴:

- 4-beams flatness optimization

Overlapping of 4 suitably appertured Gaussian beam
at a given flatness of 2% or 0.2%

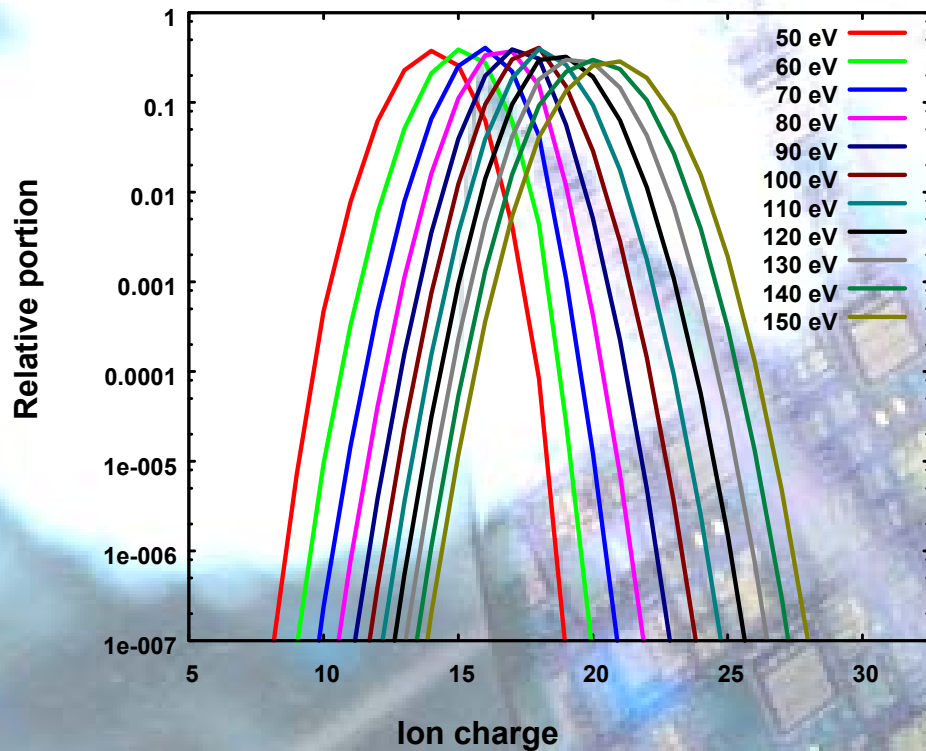


An efficiency with flatness of 0.2% is of **22%**.

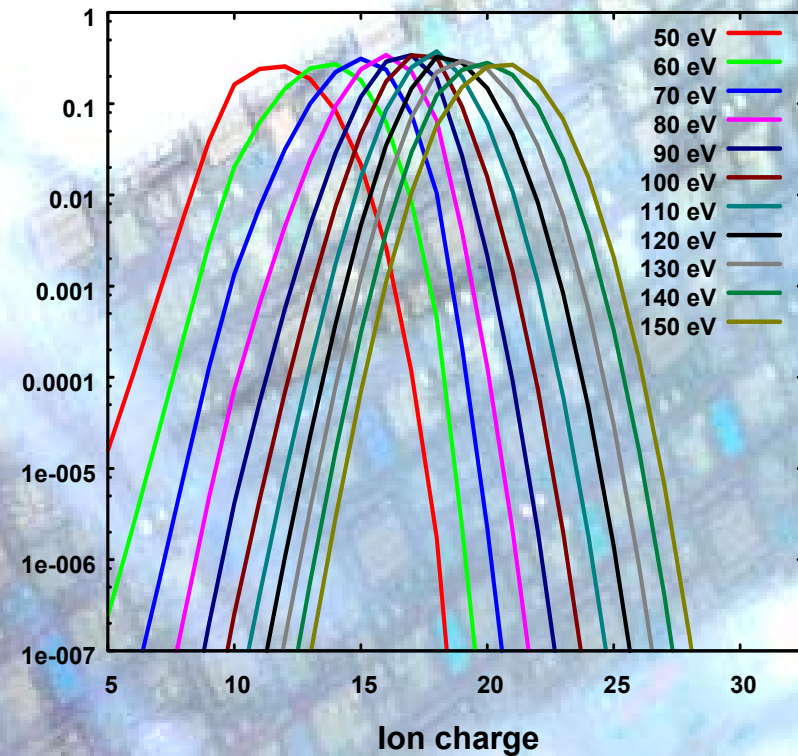
Gadolinium Plasma Emitting at 6.x nm

- Ion populations

$$N_e = 10^{19} \text{ cm}^{-3}$$



$$N_e = 5 \times 10^{20} \text{ cm}^{-3}$$



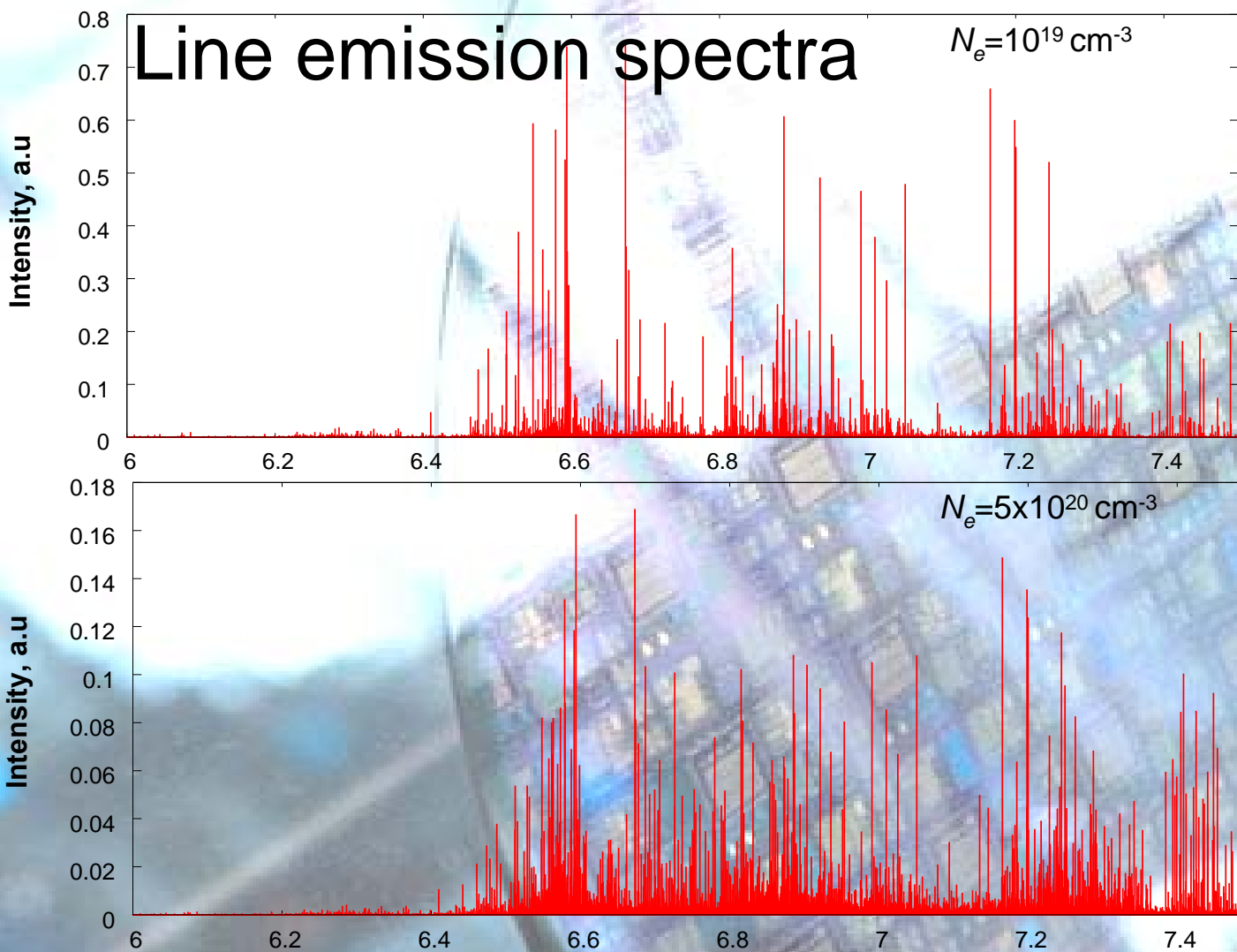
*Ion distribution spreads and average charge drops as density increases →
→ for LPP very high temperature may be necessary*

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Gadolinium Emission

low temperature regime



$T_e = 60 \text{ eV}$

$N_e = 10^{19} \text{ cm}^{-3},$
 $5 \times 10^{20} \text{ cm}^{-3}$

*$Gd^{10+} - Gd^{18+}$
are taken into
account*

*Almost 1 million
transitions in total*

*More intensive
emission is from
4f-4d transitions
($4d^9 4f^m -$
 $4d^{10} 4f^{m-1}$)*

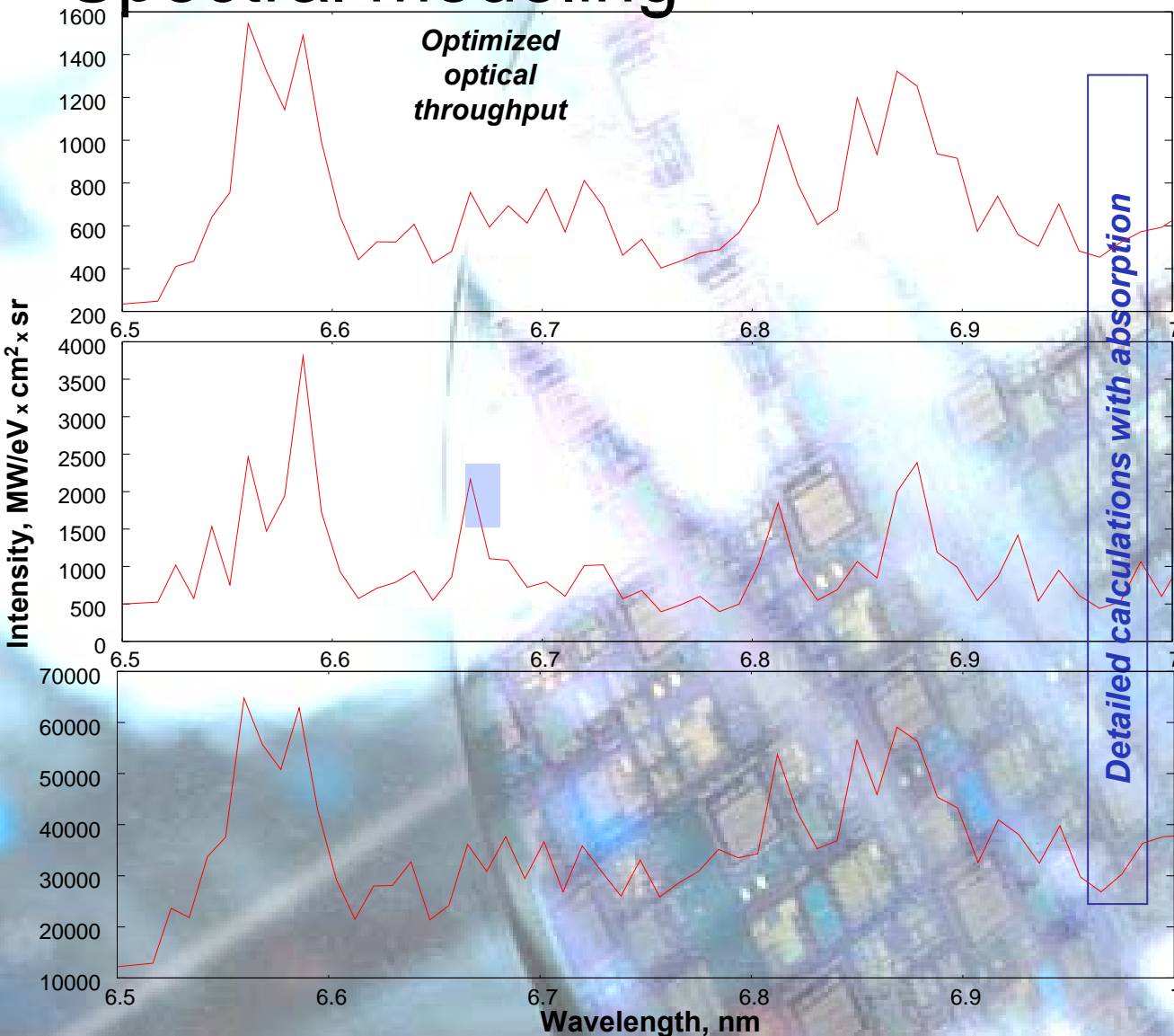
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Efficiency in Non-equilibrium Gd Plasma

low temperature regime

Spectral modeling



400 micron spherical Gd target

$N_e = 10^{19} \text{ cm}^{-3}$, $T_e = 50 \text{ eV}$

SE @ 6.68 nm of 0.6% bandwidth **6.3%**

SE @ 6.68 nm of 2% bandwidth **17.5%**

$N_e = 10^{19} \text{ cm}^{-3}$, $T_e = 60 \text{ eV}$

SE @ 6.68 nm of 0.6% bandwidth **5.3%**

SE @ 6.68 nm of 2% bandwidth **18.5%**

$N_e = 5 \times 10^{20} \text{ cm}^{-3}$, $T_e = 60 \text{ eV}$

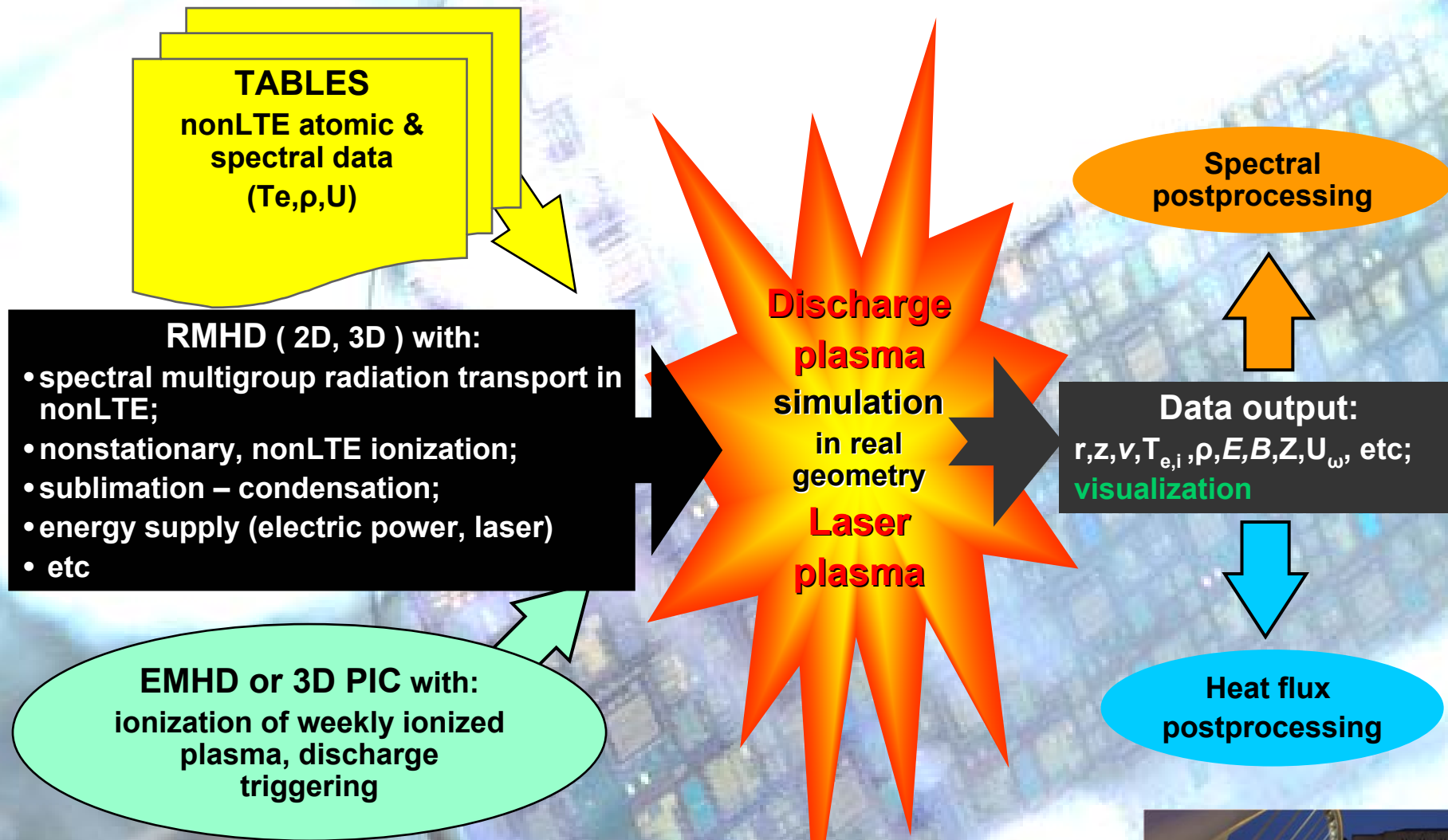
SE @ 6.68 nm of 0.6% bandwidth **5.9%**

SE @ 6.68 nm of 2% bandwidth **16.8%**

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ZETA → Z* RMHD Code → Z* BME → Z⁺ multi-physics model



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Z* Black-box Modeling Engine

Black-box Modelling Engine (Z*BME) is integrated into a specific computation environment to provide a turn-key simulation instrument, which does not require knowledge of numerical computation.

It has been adapted to simulate DPP and LPP radiation sources in a realistic geometry.

Z*BME has been installed:

in EUVA, Japan

in University College Dublin, Ireland

in Czech Technology University, Prague

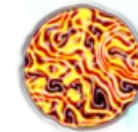


A number of joint simulations of EUV radiation sources with Z* -code of [Cymer](#), [Bochum University](#), [Xtreme Tech](#), [FOM](#), [EUVA](#), [UCD](#), [Bruker](#) has been performed in frameworks of collaborations and FACADIX, MoreMoore, Medea+, FIRE projects

2012 International Workshop
on EUV and Soft X-Ray Sources

Dublin, Ireland
October 8-11, 2012

Next Generation Modelling Tools



fire

Fluid, Ions and Radiation Ensemble
in Integrated Plasma Modelling

knowledge transfer in FP7 IAPP project

- FIRE - European FP7 Industry-Academia Partnerships and Pathways project

- The FIRE project aims to substantially redevelop the Z* code to Z+ to include improved atomic physics models and full 3-D plasma simulation of

- ✓ plasma dynamics
- ✓ spectral radiation transport
- ✓ non-equilibrium atomic kinetics with fast electrons
- ✓ transport of fast ions/electrons
- ✓ condensation, nucleation and transport nanosize particles.

- Modelling is essential in parametric scans in radiation source optimization, in fast particles and debris generation to solve current EUVL source problems as well as extending their application.



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